05.1-45 PHASE TRANSITION FROM A PARAMAGNETIC TO A MAGNETICALLY ORDERED CRYSTAL. By <u>T. Maczyński</u> and J. Kociński, Institute of Physics, Warsaw Technical University, 00-662 Warszawa, Koszykowa 75, Poland.

On the basis of Landau's theory we have calculated changes of symmetry at a continuous phase transition of a paramagnetic crystal of symmetry P6₃/mmc1' to magnetically ordered phases. The symmetry changes are induced by active representations connected with the wave vector $\mathbf{\hat{k}} = \gamma \mathbf{\hat{b}}_3$, /-0.5< $\gamma < 0.5$ /. These representations fulfil the reality condition, the Kovalev condition and the weak Lifshitz condition. We have considered two active representations which correspond to spins perpendicular to the C₆ axis. It has been proved that there can appear two types of spin structures: /1/ with constant spin length, i.e. helical structures, and /2/ with varying spin length, i.e. sinusoidal structures. For the first case, formulae for the spiral structures which are valid for any value of γ from the interval /-0.5,0.5/ have been given. We can thus obtain commensurate or incommensurate helical structures with an appropriate value of γ we can obtain the helical structure of dysprosium crystal. In Landau's theory, the corresponding phase transition can be continuous. For the second case we give an example of a sinusoidal spin structure which is characteristic for neodymium crystal. The type of magnetic group to be assigned to a particular spin structure depends on the value of γ .

05.1-46 FERROELASTIC PHASES AND ORDER PARAMETER TREATMENT IN $Pb_3(P_{1-x}As_xD_4)_2$. By <u>U.Bismayer</u>, E.Salje⁺, A.M.Glazer and J.Cosier, Clarendon Laboratory, Parks Rd., Oxford OX1 3PU, UK.

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Lead phosphate and the isostructural mixed crystals of lead phosphate-arsenate exhibit ferroelastic transitions between the phases $R\overline{3}$ m, C2/c and $P2_1/c$. High precision measurements of optical birefringence and dielectric measurements revealed at temperatures below 100K in all compounds a further transition. In lead arsenate this probably corresponds to the transformation of $Pb_3(VO_4)_2$ to the acentric phase $P2_1$.

In Pb₃(PD₄)₂ the critical order parameter exponent β fits in with a multicomponent order parameter model with q = 3, d = 2 (Salje and Devarajan, J. Phys. C (1981) 14, L1029). Along with the results of infrared reflection spectroscopy, diffuse X-ray, neutron and Raman scattering, the precursor ordering behaviour indicates dynamical fluctuations in Pb₃(PO₄)₂ and a static intermediate phase in the mixed crystals.

The transition C2/c - $P2_1/c$ is interpreted in terms of a two-order-parameter theory, which clearly demonstrates how the first-order transition in pure Pb_3 (AsO₄)₂ becomes smoother in Pb_3 ($Po_2As_{0.3}O_4$)₂.

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05.1-47 STRUCTURES OF THE FERROELECTRIC AND PARAELECTRIC PHASES OF METHYLAMMONIUM TRICHLOROMERCURATE. By M.Körfer, <u>H. Fuess</u> and J.W.Bats,Institut für Kristallographie der Universität,Frankfurt am Main,Federal Republic of Germany

A first order ferro/paraelectric phase transition in CH₃NH₃HgCl₃ at 61°C was observed by differential scanning calorimetry, by birefringence and DK-measurements. The DSC-peak exhibits considerable structure and extends over about 1°C. The ferroelectric room temperature structure has space group P32, but shows three pseudo two-fold axes (Ben Salah et al., Z.anorg.allg. Chem. (1982) <u>493</u>,178).Further refinement was achieved by the introduction of two domains linked by a two-fold axis (R(F) decreased from 0.056 to 0.046). In the paraelectric phase, the structure is described in the monoclinic space group C2 (a=13.818(2), b=7.884(1), c=9.734(3), \beta=90.49(2)). The polar three-fold axis is lost and one pseudo two-fold axis becomes a real one. Two Hg-atoms become symmetrically independent with drastically changed Hg-Cl distances as compared with the ferroelectric phase. Spectroscopic studies of the dynamics of the phase transition are in progress. The structures of the chemi-cally related compounds CH₃NH₃HgBr₃(space group P2₁/m;a=7.985(2); b=6.387(8), c=8.728(3), B = 117.17(3)) and CH₃NH₃HgJ₃ (space group P2₁/n; a=9.089(4), b=7.006(4), c=14.817(11), B=91.70(4)) were also determined, but these compounds are not ferroelectric at room temperature.

05.1-48 A NEUTRON SCATTERING STUDY OF THE INCOMMENSURATE PHASE IN ALPO4 (BERLINITE) NEAR THE*α*-*β* TRANSITION.By H. Arnold⁺, J. Bethke⁺, G. Eckold⁺, Th. Hahn⁺, J. Min-Hua⁺⁺; + Inst. für Kristallographie, RWTH Aachen, 51 Aachen, ++ Inst. für Festkörperforschung, Kernforschungsanlage Jülich, 5170 Jülich, FRG; +++ Crystal Material Institut, Shandong University, Jinan Shandong, The Peoples Rep. of China.

Quartz and its isotype ALPO4 (Berlinite) have many physical properties in common, especially the symmetry change at the α - β transition and the microdomain structure near the transition temperature. We report on our elastic neutron scattering results of the intermediate incommensurate phase of ALPO4 existing in a temperature range of about 1.5 K between the α and the B phase. The experiments were performed on the triple axis spectrometer UNIDAS at the FRJ-2 reactor in Jülich using 0.25° collimation. Crystals with typical dimensions of 2.5 x 2.5 x 20 mm³ were grown under hydrothermal conditions. After having passed the α - β transition several times, satellite reflections were observed in the directions a*, b* and a* - b* in various Brillouin zones. In particular, satellites correspondent ponding to directions of the modulation vector $\overset{q}{}$ perpendicular to the reciprocal lattice vector \underline{Q} exhibit large structure factors. The temperature dependence of q, which indicates the incommensurate character of this intermediate phase has been determined unambigously. Within the temperature interval of 1.5 K the satel-lites display a linear shift from q = 0.02 up to q = 0.038 reciprocal units with increasing temperature.