## 10-Physical and Chemical Properties of Materials in Relation to Structure (Superconductors, Fullerenes, etc)

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25% are less favourable. The second orientation of  $C_{60}$ found in the low-temperature structure could not be identified at room temperature.

The orientational disorder of C<sub>60</sub> in the structure of (hydroquinone)<sub>3</sub>C<sub>60</sub> (trigonal, R3m) has been investigated at four temperatures, 100K, 200K, 293K and 373K. C<sub>60</sub> occupies clathrate-like cages with symmetry 3m (Ermer, O., Helv. Chim. Acta 1991, 74, 1339-1351). At 100K,  $C_{60}$  shows a threefold orientational disorder, the crystallographic molecular symmetry being 2/m. At higher temperatures, disorder increases: at 373K, it is represented by two sets of oriented symmetry-related molecules with site symmetry 1. One is near the orientation found at 100K, the other places a molecular threefold axis close to the crystallographic threefold axis. As for pure C60, this parametrization permits to determine a probability density function of molecular reorientation.

MS–10.02.03 Defects and phase transitions in  $\mathrm{C}_{70}$  and  $\mathrm{C}_{60}.$ 

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Two subjects will be addressed in some detail. (i) Phase transitions in C70 [ G. Van Tendeloo et al. Europhys. Let.21 (1993) 329]. Crystals grown from the vapour under reduced helium pressure exhibit two different morphologies:

the majority is cubic, a few % only is hexagonal. Electron diffraction on fragments obtained by crushing under liquid nitrogen shows that at room temperature the crystal structures are respectively predominantly FCC and HCP with a c/a ratio close to ideal consistent with their morphology.

In most cases the structure is faulted.

At liquid nitrogen temperature the HCP crystals exhibit an orientationally ordered phase with a superstructure which can be described as monoclinic with  $a_m = a_H$ ,  $c_m = 2a_H$ ,  $b_m = c_H$ and  $\beta$  close to 120°. This low temperature phase can occur in six orientation variants with respect to the hexagonal room temperature phase. The layers of molecules are stacked according to the HCP mode (c/a  $\equiv$  1.82) with their long molecular axis parallel to  $b_m$ . Rows parallel to  $a_m$  consist of parallel molecules having the same azimuthal orientation, but in adjacent rows the molecules differ 180° in orientation. A second orientational stacking mode of the same type of layers leading to  $b_m = 2c_H$  is occasionally found (orientational polytypism). With increasing temperature the different orientational degrees of freedom of the molecules are gradually excited. In a first transition (~ 276 K) the orientational order dissappears as a consequence of the free rotation about the long molecular axis; hereby c/a remaining 1.82. At 337 K c/a decreases to 1.66; subsequently the ABAB stacking is

irreversibly changed into a rhombohedral ABCABC stacking by a shear transformation; finally, at higher temperatures the structure becomes FCC.

(ii) The defect structure of C60 [S. Muto, G. Van Tendeloo, S. Amelinckx, Phil.Mag (1993) in the pressl.

The orientational disorder-order transition from face centered cubic into simple cubic causes fragmentation of single crystals into domains related by the 90° symmetry rotation, lost in the phase transition; also translation variants are produced. This domain structure has been studied by diffraction contrast electron microscopy. A further transition from this primitive cubic structure into one with doubled cell parameters, termed "2ao" superstructure has been observed by electron diffraction.

All the defects which commonly occur in low stacking fault energy face centered cubic alloys are also found in C<sub>60</sub> at room temperature. High resolution electron microscopy has allowed to image these defects on a molecular scale and we will discuss some of these defects in detail:

Twin interfaces and the lattice relaxation along them.

- (ii) Intrinsic and extrinsic stacking faults and their relation with partial dislocations.
- (iii) Lines of partial vacancies along the intersection lines of stacking faults.
- (iv) Interaction of stacking faults and twin boundaries.
- (v) Surface reconstruction of C<sub>60</sub>, probably because of the presence of oxygen at the surface.

## MS-10.02.04 PHASE TRANSFORMATIONS OF C60 UNDER NONHYDROSTATIC PRESSURE

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Fullerenes under pressure present several new phenomena depending on pressure and temperature conditions. Under hydrostatic compression, C<sub>60</sub> molecules are stable up to 20GPa (Duclos *et al.* Nature, 1991, 351, 380). However, application of nonhydrostatic pressure above 15 GPa transforms fullerenes into polycrystalline diamond at room temperature, the color of the bulk sample being yellowish by transparency (Nûnez Regueiro *et al.* Nature, 1992, 355, 237). Furthermore, application of nonhydrostatic pressure at high 237). Furthermore, application of nonhydrostatic pressure at high temperature give rise to black non insulating samples. We report on the x ray synchrotron radiation and electron diffraction characterization of these transformations.

Electron diffraction patterns of samples transformed at room temperature present continuous rings corresponding to all f.c.c. diamond reflections or diffuse rings of an amorphous phase. The diffraction patterns evolve continuously from diamond rings to amorphous rings and the sample morphology, as observed by optical and electron microscopy, is not lamelar. From the dark-field images, we deduce that the size of the crystallites varies from 20 to 1000A in the different chips. In all samples the amorphous phase exists, the largest proportion of well crystallized diamond material was found when the applied pressure was the least hydrostatic. We have analyzed the diffuse pattern of samples which seemed totally amorphous by using x ray synchrotron radiation at DCI (LURE) It corresponds to using x ray synchrotron radiation at DCI (LURE). It corresponds to  $sp^3$ -type carbon and is always mixed with a small amount of polycristallized f.c.c. diamond. Several simulations and fits show that this  $sp^3$  amorphous phase contains, together with staggered C-C bonds, some amount of eclipsed C-C bonds as found in Lonsdaleite. Simulation spectra of theoritical phases containing C<sub>20</sub> or C<sub>24</sub> polyhedra do not agree with our measurements and the number of eclipsed C-C bonds must be limited.

On the contrary, samples obtained in a belt press at moderate pressures and high temperatures are black non insulating and amorphous. Electron diffraction and x ray synchrotron radiation characterization of this phase show its  $sp^2$ -character, affected by preferential orientation.