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Microstructures in Amorphous Nb₃Ge Films

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In amorphous Nb_3Ge films made by DC getter sputtering, small-angle X-ray scattering indicated that amorphous oxide clusters were present.

Superconducting Nb₃Ge films (Gavaler, 1973; Testardi, Wernick & Royer, 1974) were prepared by DC getter sputtering onto a cold (300 K), amorphous SiO₂ substrate. Films, approximately 1 μ m thick, were carefully examined by largeangle X-ray scattering using a Huber camera in the Seeman-Bohlin geometry. The X-ray diffraction patterns exhibited three broad peaks and no crystalline reflections (Chencinski & Cadieu, 1974). The intensity of small-angle X-ray scattering (SAXS) from these amorphous Nb₃Ge films on SiO₂ substrate were measured for 0.01 $\text{\AA}^{-1} \leq |\mathbf{K}| \leq 0.1 \text{\AA}^{-1}$, $(|\mathbf{K}| =$ $4\pi \sin \theta/\lambda$ with Cu K radiation and a Kratky camera. With proper corrections for substrate scattering and absorption (Chi & Cargill, 1976), the radius of gyration for the scattering regions was obtained from a Guinier plot and found to be 71 Å. The radii of gyration for samples, which were heat treated for 5h at 550, 620, 690 and 735°C in a vacuum system (10⁻⁶ torr), were 83, 85, 88 and 127° Å respectively. The large-angle diffraction patterns for samples annealed at 690° and 735°C showed crystalline reflection lines. In both cases, detailed examination of the crystalline pattern clearly indicated that most of the sample was of Nb₅Ge₃ phase. Some reflection lines may also be attributed to A15 Nb₃Ge, GeO₂ and NbO. The X-ray patterns for crystallized samples and Rutherford back-scattering measurements confirmed that there was oxygen in the sample. These observations indicated that amorphous oxide clusters were the major microstructures in these amorphous Nb₃Ge films.

The full paper will be submitted to *Materials Science and Engineering*.

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Small-Angle X-ray Scattering on Ferromagnetic Colloids

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The size distribution and the structure of colloidal cobalt particles in toluene, as well as interparticle correlations, have been studied.

Ferrofluids are colloidal suspensions of ferromagnetic monodomain grains in thermal equilibrium with a nonmagnetic and insulating fluid. We are using magnetic colloids of cobalt in toluene obtained from the thermal decomposition of cobalt octacarbonyl in the presence of a terpolymer dissolved in toluene (Liebert, Martinet & Strzelecki, 1972).

A quantitative analysis of electron micrographs shows that the radius distribution of cobalt spheres corresponds to a log-normal law with a mean radius \overline{R} of about 52 Å.

The experimental magnetization curve may be described by Langevin curves weighted by a coefficient proportional to the log-normal law, F(R), (Anthore, Petipas, Chandesris & Martinet, 1977): $\overline{R}_m = 46$ Å; $\sigma = 1.28$. \overline{R}_m is always smaller than \overline{R} . Size distributions and structure of particles: Let $j_n(S)$ be the normalized function. The $\log j_n(s)$ versus s^2 curve shows a good linearity, which is characteristic of a small variation in the particle size, but the experimental curves $s^3 j_n(s)$ have a non-standard form. They tend towards a smaller limiting value than that which would be expected in connection with the discontinuity between pure cobalt and toluene; again it decreases with the aging time at constant radius of gyration. The experimental values are given in Table 1.

The volume fraction β occupied by the particles can be calculated by different methods: β_{abs} is the value determined from the absorption measurement of the incident beam, β_{Q0} that from the integrated intensity ($\beta_{Q0} = 4.2 \times 10^{-3}$), β_G that from the intensity $j_n(0)$ ($\beta_G = 4.1 \times 10^{-3}$), and β_M that from the magnetization at saturation; β_{ch} is calculated from the quantity of cobalt octacarbonyl decomposed ($\beta_{ch} = 4.6 \times 10^{-3}$). The difference between β_M and the other values reveals that a fraction of cobalt cannot be considered as free