contribution we introduce a program that is aimed to be a swiss army knife for processing 3D-PDF and corresponding diffuse scattering data. Refinement tools are not included in this program and will be discussed elsewhere [4].

The features of the program currently include:

- quick and easy navigation through large 3D volumes,
- voxel precise editing of 3D data sets that may contain even more than one billion data points,
- separation of Bragg, diffuse and background scattering,
- symmetry averaging,
- transformation between reciprocal space and 3D-PDF space,
- tools that allow establishing a disorder model from 3D-PDF data.

To fulfill the requirements visualization of the data is done by pixel based imaging of 2D cuts through 3D data volumes instead of using iso-surface representations. Currently, the program supports its own format and the XCAVATE [5] file format, but other formats will be included if required. The program is by no means restricted to visualization and manipulation of 3D-PDF data, but may be used as a general editor for diffuse scattering data.

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Keywords: pair distribution function, diffuse scattering, data processing software

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Computer Simulation of Diffuse Scattering in Fe(II) Spin Crossover Compounds. M.Zubko^a, R.Neder^b,

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Many Fe (II) spin crossover compounds show a thermal transition from the ${}^{1}A_{1g}$ low spin (LS) ground state to the excited ${}^{5}T_{2g}$ high spin (HS) state. As a consequence the metalligand bond lengths increase by up to 0.2Å [1]. This increase has a strong impact on the whole crystal lattice as observed from the large temperature variation of the lattice parameters in the spin crossover region.

The compound $[Fe(ptz)_6](BF_4)_2$ (ptz=1-propyltetrazole) belongs to the group of the octahedrally coordinated Fe(II) spin crossover compounds. It crystallizes in space group R-3 (Z = 3) and the structure consists of exactly trigonal, neutral layers perpendicular to the *c* axis [2,3]. The iron atoms occupy special position 3(a) and all six ligands are equivalent. The (BF₄)⁻ anions are placed on the threefold axis.

Slow cooling through $T_{1/2} = 123K$ causes appearance of diffuse scattering along the c^* direction and splits Bragg peaks into two maxima. This indicates the presence of short range order in the low temperature diffuse phase [4].

Simulations and refinement of the disordered phase have been done with the DISCUS program package [5,6]. The simplest applied model consist of stacks of identical layers, which are stacked by either the perfect [$\frac{1}{3}$, $-\frac{1}{3}$, $\frac{1}{3}$] vector or with an additional shift of [$\frac{1}{3}+\delta_x$, $-\frac{1}{3}+\delta_y$, $\frac{1}{3}$]. The model is characterized by three parameters describing: the additional shift between neighbouring layers, the percentage of additionally shifted layers and the correlations along the *c* axis between shifted layers. The refinement of the parameters was done through the differential evolutionary algorithm implemented in the DISCUS software package.

Computer simulations of X-ray diffuse scattering give insight into the real structure of the disordered phase of $[Fe(ptz)_6](BF_4)_2$ compound. The change of the peak shape with time and temperature can be explained by formation of lamellar domains perpendicular to the *c* direction. Such deformation becomes possible due to decrease of the interlayer Van der Waals forces caused by spin-crossover transition.

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The Debye scattering formula in *n* dimensions Thomas Wieder

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Consider an Euclidian coordinate system E_n of dimension n = 1,2,3,... Let $r_{i,j}$ denote the interatomic distance vector among atoms *i* and *j* in an atomic assembly of *M* atoms within E_n . The contribution of $r_{i,j}$ to a diffraction pattern is obtained from an integration over the *n*-dimensional sphere of radius $|r_{i,j}|=r_{i,j}$. One arrives at the Debye scattering formula I(k,n) valid in E_n . With *k* as the length of the scattering vector and with f(k,i) as the atomic form factor the *n*-dimensional Debye scattering formula is

$$I(k,n) = \sum_{i=1}^{M} \sum_{j=1}^{M} f(k,i) f^{*}(k,j) F(k,n)$$
(1)

where one has

$$F(k,n) = \frac{2^{\frac{n}{2}-1}}{2^{\frac{n}{2}-1}} \Gamma\left(\frac{n}{2}\right) (kr_{i,j})^{-\frac{n}{2}} \left(J\left(\frac{n}{2}, kr_{i,j}\right)n - J\left(\frac{n}{2}+1, kr_{i,j}\right)kr_{i,j}\right)$$

with $\Gamma(x)$ as the Gamma function and J(l,x) as the *l*-th Bessel function of the first kind.

For a given dimension *n* the formula (1) takes on a concrete form and such forms will be given for n=2,...,8.

However, an application for cases with n > 3 is not known yet.

Keywords: scattering theory, amorphous scattering, analysis of disordered structures