

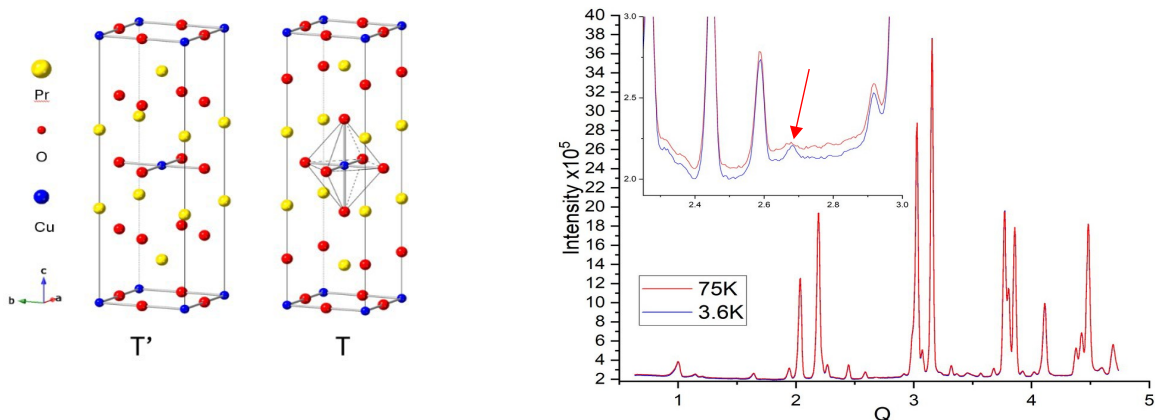
## Poster

Restoring superconductivity in hole-doped  $\text{Pr}_2\text{CuO}_4$  via high-pressure phase transitionsM. Edwards<sup>1</sup>, M. S. Senn<sup>1</sup><sup>1</sup>Department of Chemistry, University of Warwick, Coventry, CV4 7AL, United Kingdom matt.edwards@warwick.ac.uk

The hole-doped lanthanum cuprates, based on Ruddlesden-Popper  $n = 1$  (RP1)  $\text{La}_2\text{CuO}_4$ , have attracted intense focus since the discovery of high-temperature superconductivity in  $\text{La}_{2-x}\text{Ba}_x\text{CuO}_4$ . The temperature vs doping level ( $x$ ) phase diagram for these materials includes a region around  $x = 0.125$  (1/8th) featuring a marked drop in superconducting critical temperature ( $T_c$ ). This region is understood to coincide with both the emergence of charge density wave ordering, which seemingly competes with superconductivity, and a structural phase coexistence between the low temperature orthorhombic (LTO) and low temperature tetragonal (LTT) structures which arise from cooperative tilts of the  $\text{CuO}_6$  octahedra present in the RP1 structure about the  $[110]$  (LTO) or alternating between the  $[100]/[010]$  axis (LTT). This structural phase coexistence is not observed in systems where the A-site cation size variance ( $\sigma_A$ ) is low, e.g.  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  [1], leading us to believe that  $\sigma_A$  acts as a chemical control parameter for tuning the phase separation. Thus, we have devised a model system,  $\text{Pr}_{2-x}\text{Ca}_x\text{CuO}_4$ , where  $\sigma_A$  is zero within the precision of the reported Shannon ionic radii, which should prove a promising candidate for gaining insights into the interplay of electronic properties and structure in these superconducting materials.

$\text{Pr}_{2-x}\text{Ca}_x\text{CuO}_4$  crystallises in the  $\text{Nd}_2\text{CuO}_4$  type structure ( $T'$ ), which differs from the  $\text{K}_2\text{NiF}_4$  RP1 structure ( $T$ ) adopted by the lanthanum cuprates most obviously in the lack of apical oxygen, leading to differing coordination of the Cu – four-fold in the  $T'$  and six-fold in the  $T$  (Fig. 1(a)) [2]. We have synthesised  $\text{Pr}_{1.875}\text{Ca}_{0.125}\text{CuO}_4$  and shown by resistivity and susceptibility measurements that it is not superconducting but rather an antiferromagnetic insulator, similar to the undoped parent compound  $\text{Pr}_2\text{CuO}_4$ . This is not surprising as, typically, superconductivity is only induced in cuprates with the  $T'$  structure by electron doping, i.e. with  $\text{Ce}^{4+}$ . However, there are indications in the literature that  $\text{Pr}_2\text{CuO}_4$  undergoes a high-pressure phase transition to the  $T$  phase at about 15 GPa [3] which, if we can replicate in our Ca doped system, should allow us to convert it to a structural symmetry where superconductivity can be stabilised.

Initial attempts to stabilise the  $T$  phase in our  $\text{Pr}_{1.875}\text{Ca}_{0.125}\text{CuO}_4$  sample at high-pressure via *in situ* PXRD experiments have been undertaken. Using diamond anvil cells, we have measured up to 7 GPa with no transition observed, indicating higher pressures are required. Despite this there are several interesting features associated with the  $T'$  phase including neutron diffraction data showing the emergence of a single magnetic peak at low temperatures (Fig. 1(b)) which does not appear to have the same  $k$ -vector as the parent material which suggests the magnetic ordering has changed. High-resolution PXRD data shows increasing peak broadening with temperature suggesting unusual microstrain behaviour which could be indicating the onset of a phase transition.



**Figure 1.** (a): schematics of the  $T'$  and  $T$  structures. (b): WOMBAT neutron data with inset highlighting the emerging magnetic peak.

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