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# Characteristics and local structure of hafnia-silicate-zirconate ceramic nanomixtures

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Zirconate systems having the composition  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$ , where x = 2, 7 and 12 mol%  $Y_2O_3$ , were synthesized by a sol-gel method. The analysis of X-ray diffraction data showed the presence of the t-ZrO<sub>2</sub>, m-ZrO<sub>2</sub>, m-HfO<sub>2</sub>,  $Y_2SiO_5$  and  $Y_2Si_2O_7$  crystalline phases in a ceramic nanomixture. Spectroscopic data show that the increase of the  $Y_2O_3$  content of samples determines the increase of the t-ZrO<sub>2</sub>, m-HfO<sub>2</sub> and silicate crystalline phases. Gap energy values decrease almost linearly with increasing  $Y_2O_3$  content of samples. A detailed study of XANES data does not show a significant difference with increasing  $Y_2O_3$  content of the samples suggesting an appreciable stability of the hafnium ions +4 oxidation state and their microvicinity. EXAFS results show that the local structure around the Hf cation is similar to that from the monoclinic crystalline HfO<sub>2</sub> where the Hf–O coordination number tends to 7. The bond lengths of Hf–O shells show small deviations from ~2.12 Å and the Hf–metal paths become more structured by increasing the  $Y_2O_3$  content of the samples.

#### 1. Introduction

Zirconium oxide based systems have been drawing increasing attention of researchers due to their unique properties that are promising for biomedical applications. Thus, ceramics based on  $ZrO_2$  are well known for their low thermal conductivity, high fracture toughness and biocompatibility in comparison with other materials (Ponnilavan *et al.*, 2015).

As has been observed up to now, pure zirconia  $(ZrO_2)$ possesses three polymorph structures: cubic (c-ZrO<sub>2</sub>), tetragonal (t-ZrO<sub>2</sub>) and monoclinic (m-ZrO<sub>2</sub>). The t-ZrO<sub>2</sub> and m-ZrO<sub>2</sub> phase transformation has important technological consequences, since the extensive microcracking caused by the anisotropic volume expansion accompanying the reaction leads to the loss of mechanical integrity in sintered bodies. When a metal oxide with a lower valence cation is incorporated into cubic ZrO<sub>2</sub>, the fluorite-type cubic phase is stabilized. One of the common 'stabilizers' is yttrium oxide since it has good physical and chemical stability (Torres & Llopis, 2009; Qin et al., 2012; Singh et al., 2015). The substitution of this lower valence cation (+3) in the  $Zr^{+4}$  sublattice sites creates oxygen vacancies. The Y<sub>2</sub>O<sub>3</sub> can stabilize the fluoritetype cubic structure and hence avoid the problem of the t-ZrO<sub>2</sub> and m-ZrO<sub>2</sub> transformation by eliminating the t-phase (Wu et al., 1992; Chu et al., 2016).

Further, the dual role of zirconia in glass matrices was confirmed as a network former and/or as a network modifier. Thus, zirconia playing the network former role leads to a network consisting of  $ZrO_4$  tetrahedra while playing the

network modifier role leads to a network consisting of  $ZrO_6$  octahedra with four bridging and two non-bridging oxygen atoms (Glazneva *et al.*, 2012). Addition of HfO<sub>2</sub> to  $ZrO_2$  enhances the chemical stability and thermal-shock resistance of the host ceramic (Zhen *et al.*, 2007; Karunarathne *et al.*, 2017). Note that zirconia containing silicon oxide has attracted attention too because of its excellent resistance to alkali corrosion and its low thermal expansion (Simhan, 1983; Kwon *et al.*, 2017). In such a material, hafnium oxide or zirconium oxide were the best metal oxides to replace the silicon dioxide.

Note that fundamental studies on the crystalline structure and microstructure of monoclinic  $ZrO_2$  ceramics with or without dopants have rarely been reported (Kwon *et al.*, 2017; Rada, Culea *et al.*, 2018). The research on monoclinic  $ZrO_2$ ceramics can be associated with difficulties in the manufacturing process and the mechanical instability at temperature variations.

The  $ZrO_2-Y_2O_3$  compositions are candidates for the transmutation of radionuclides in reactors and as acceleratorbased neutron sources due to their high radiation resistance (Moll *et al.*, 2009).

Taking into account the important properties of the oxides presented above, we prepared some new ceramics containing all of these oxides using the sol-gel method. Considering the development of new materials with mixed compositions, the experimental conditions can be adjusted in order to obtain stable or metastable compounds. The physical and chemical properties of mixed oxides prepared by the sol-gel method are of great interest due to their thermal and chemical properties and because they are highly acidic due to the presence of OH groups on the surface.

Bearing in mind the above-mentioned aspects, in this study we wanted to prepare some new zirconate materials with potential biomedical and optical applications. Thus, new zirconate nanostructures were prepared in the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system where x = 2, 7 and  $12 \text{ mol}\% Y_2O_3$  via the sol-gel method. To the best of our knowledge, this system has not been reported in the scientific literature until now.

The investigation of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$ zirconate materials was carried out by X-ray diffraction (XRD) and scanning electron microscope (SEM) analysis, Fourier transform infrared (FT-IR), UV-Vis, photoluminescence (PL), Raman and X-ray absorption spectroscopies. X-ray absorption spectroscopy (XAS) with its two amendments, X-ray absorption near-edge structure (XANES) and extended X-ray absorption fine structure (EXAFS), represents an important investigation technique of the local structure around ions of interest. Thus, EXAFS analysis was used to study the structure of the Hf atoms' microvicinity, their coordination number and the length of the bonds involving Hf atoms.

First, one of the main objectives of the present work was to perform the structural characterization of the studied new zirconate materials. Our interest was focused on obtaining structural information (via XRD, SEM analysis, FTIR and X-ray absorption spectroscopies) concerning the characteristic crystalline phases and structural units that occur in the studied materials. Second, some spectroscopic properties of the studied zirconate samples were investigated (by using FT-IR, UV-Vis, PL and Raman spectroscopies) to find out their potential applications for optical devices. The electronic transitions that take place were studied. Finally, the influence of the addition of  $Y_2O_3$  was followed, since dopants can be useful tools to improve the structural and/or behavioral properties of the host materials.

#### 2. Experimental data

Tetraethoxysilane and high purity ZrO2, Y2O3 and HfO2 powders were used as starting materials. Tetraethoxysilane was diluted with ethanol. Solutions of Zr(NO<sub>3</sub>)<sub>4</sub>, Y(NO<sub>3</sub>)<sub>3</sub> and Hf(NO<sub>3</sub>)<sub>4</sub> were prepared by dissolving ZrO<sub>2</sub>, Y<sub>2</sub>O<sub>3</sub> and HfO<sub>2</sub> into  $HNO_3$  solutions of pH = 2. Suitable volumes of tetraethoxysilane, Zr(NO<sub>3</sub>)<sub>4</sub>, Y(NO<sub>3</sub>)<sub>3</sub> and Hf(NO<sub>3</sub>)<sub>4</sub> were mixed and stirred at 70°C to obtain homogeneous solutions. The resulting solutions were stirred at 100°C until a dry powder was obtained after the evaporation of  $H_2O$ . These dried powders were uniaxially compacted under pressure in the form of disks (2 mm thick, 9 mm diameter). The disks were calcined in alumina crucibles using an electric furnace heated at 1400°C for two hours. The disks were put into the electric furnace direct at this temperature. After two hours, the melted disk materials were quenched at room temperature by pouring onto a stainless-steel plate.

The samples were characterized by several methods: XRD, SEM analysis, FTIR, UV-Vis, PL, Raman and XAS spectroscopy that were previously described by Rada *et al.* (2013, 2017, 2019), Rada, Zhang *et al.* (2018) and Dehelean *et al.* (2016).

The changes of the microstructure of the samples by doping were characterized by analysing the profile of the diffraction peaks and SEM micrographs. FTIR spectroscopy was recorded to determine the presence of the specific structural units, such as zirconia or silicate network in the prepared samples. UV-Vis data were used to determine the gap energy. The UV-Vis, PL and Raman spectroscopy were also used to show the presence of the zirconia polymorphs. The XRD peaks of monoclinic zirconia and monoclinic hafnia are very similar because they arise at nearly the same angles. XANES and EXAFS were performed in order to indicate the formation of the HfO<sub>2</sub> crystalline phase with monoclinic structure.

#### 3. Results and discussion

#### 3.1. XRD data

X-ray diffraction was used to identify the main crystalline phases in the studied samples. Fig. 1 shows the XRD patterns of the samples from the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system where x = 2, 7 and  $12 \text{ mol}\% Y_2O_3$ . All diffractograms show peaks characteristic of crystalline phases. In order to identify the crystalline phases that appear in the diffraction



X-ray diffraction patterns for the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system.

patterns we followed a standard procedure, namely computer analysis based on the powder diffraction file (PDF).

This analysis permitted the identification of five crystalline phases present in the samples. Two of these phases belong to the ZrO<sub>2</sub> crystalline phase with monoclinic (PDF No. 03-065-1024, monoclinic lattice with a = 5.14608 Å, b = 5.21177 Å, c =5.31302 Å and  $\alpha = 99.222^{\circ}$  cell parameter) and tetragonal (PDF No. 00-050-1089, a = 3.5984 Å, c = 5.152 Å) symmetry. The HfO<sub>2</sub> crystalline phase with monoclinic structure (PDF No. 00-034-0104, monoclinic lattice with a = 5.2851, b = 5.1819, c = 5.1157 Å cell parameter), the yttrium silicate (Y<sub>2</sub>SiO<sub>5</sub>, PDF No. 00-021-1456) and yttrium disilicate (Y<sub>2</sub>Si<sub>2</sub>O<sub>7</sub>, PDF No. 00-021-1459) crystalline phases were also shown in the studied samples. The assignments of the most important crystalline peaks are shown in Fig. 1. The X-ray diffractograms from Fig. 1 show that by increasing the yttrium content of samples, *x*, the relative intensity of the peaks attributed to the monoclinic ZrO<sub>2</sub> decreases, while the relative intensity of peaks assigned to the tetragonal ZrO<sub>2</sub> phase increase. Bearing in mind different potential applications, this is an important evolution since the tetragonal ZrO<sub>2</sub> is more stable that the monoclinic one.

The (111) and (111) reflection peaks situated at about 28.2° and 31.4° of the monoclinic  $ZrO_2$  crystalline phase (Rada, Culea *et al.*, 2018) did not appear to be significantly modified by increasing the Y<sub>2</sub>O<sub>3</sub> content of the samples. Analysis of the reflections in the XRD patterns of the monoclinic  $ZrO_2$  phase and monoclinic HfO<sub>2</sub> crystalline phase shows the same satellite peaks. Accordingly, it is difficult to follow the effect of the Y<sub>2</sub>O<sub>3</sub> content on the amount of monoclinic  $ZrO_2$  phase in the studied samples.

A small amount of  $Y_2Si_2O_5$  crystalline phase was detected in the sample with  $x = 2 \mod \% Y_2O_3$ . By increasing the  $Y_2O_3$ content up to  $x = 7 \mod \%$ , the intensity of the characteristic peaks of the  $Y_2SiO_5$  crystalline phase increases. After that, at a higher  $Y_2O_3$  content, x = 12 mol%, a new diffraction peak centered at about 22.03° can be seen suggesting the presence of the  $Y_2Si_2O_7$  crystalline phase. This compositional evolution shows that by increasing the yttrium ions content of samples the zirconia host matrix was converted into a yttrium–silicate– zirconia matrix. The excess of oxygen atoms can be accommodated in the host matrix by the formation of yttrium silicates.

#### 3.2. FT-IR data

The FT-IR spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  samples with x = 2, 7 and  $12 \mod \%$  are shown in Fig. 2. Each FT-IR absorption band is related to some vibration type specific of the structural groups that build up the investigated samples. The positions and assignments of the most important absorption bands shown in the FT-IR spectra are discussed. The assignments were made based on data from the literature concerning some vitreous or crystalline compounds related to the system investigated in this work (Lee & Condrate, 1988; Choi & Choy, 2016; Vitanov *et al.*, 2014; Ortiz *et al.*, 2005; Khomenkova *et al.*, 2014; Najafinezhad *et al.*, 2017; Ferdov *et al.*, 2012).

According to Lee & Condrate (1988), the wavenumber and/ or intensity of vibrational modes involving Zr-O bonds will vary depending upon the arrangement of the oxygen atoms around the zirconium atoms. Thus, several zirconium–oxygen structural units can be considered with respect to the Zr-Ostretching modes:  $ZrO_4$  units,  $ZrO_6$  units, monoclinic zirconiatype  $ZrO_7$  units, tetragonal zirconia-type  $ZrO_8$  units and cubic zirconia-type  $ZrO_8$  units.

The absorption bands around 436, 386 and  $310 \text{ cm}^{-1}$  are present in all the studied samples (Fig. 2). These bands can be assigned to the symmetric and/or asymmetric stretching vibrations of chemical bonds in which zirconium and hafnium ions are involved.

The FT-IR features situated at about 525 cm<sup>-1</sup>, 580 cm<sup>-1</sup> and 745 cm<sup>-1</sup> are due to Zr–O–Zr asymmetric and Zr–O



Figure 2 Infrared spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system. stretching modes and confirm the presence of monoclinic  $ZrO_2$  phases in the studied system. The intensity of the mentioned vibrational features becomes less prominent with increasing the Y<sub>2</sub>O<sub>3</sub> content of samples. In addition, a new weak IR feature associated with the Zr–O vibrations in tetragonal ZrO<sub>2</sub> crystalline structure appears around 600 cm<sup>-1</sup> for x = 12 mol% Y<sub>2</sub>O<sub>3</sub>. Note that this crystalline phase was detected by XRD analysis, too.

The IR bands centered at about 870, 960 and 1020 cm<sup>-1</sup> can be associated with the symmetric and antisymmetric stretching vibrations of the silicate units. By increasing the  $Y_2O_3$  dopant level up to 12 mol%, the intensities of these bands increase, suggesting the formation of higher amounts of yttrium silicates in the host matrix.

The structural changes observed by the FT-IR investigation suggest that the yttrium ions play a network modifier role in the studied  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system. These structural changes, determined by increasing the Y<sub>2</sub>O<sub>3</sub> content of the samples, consist of the increase of the tetragonal ZrO<sub>2</sub> and silicate crystalline phases. This assumption is in good agreement with the results of the XRD investigation. The observed compositional evolution shows that the Y<sub>2</sub>O<sub>3</sub> content can be used as an efficient tool to improve the stability of the studied system.

#### 3.3. UV-Vis data

The UV-Vis spectra of the  $3\text{HfO}_2 \cdot 15\text{SiO}_2 \cdot xY_2\text{O}_3 \cdot (82 - x)\text{ZrO}_2$  zirconate nanostructures where x = 2, 7 and 12 mol% are given in Fig. 3. The samples show a very intense absorption in the UV region and a moderate absorption in the visible region. The bands from 225 nm, 265 nm and 350 nm correspond to the metal-ligand charge-transfer (oxygen electron from the free orbital of the metal ion). The absorption band from 225 nm is the highest one and is assigned to the charge transfers (i) from the oxygen anions to the zirconium cations  $(O^{2-} \rightarrow Zr^{4+})$  in ZrO<sub>2</sub> and (ii) from the oxygen anions to the



**Figure 3** UV-Vis spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system.

hafnium cations ( $O^{2-} \rightarrow Hf^{4+}$ ) in HfO<sub>2</sub> (Fan *et al.*, 2017; Qi & Zhou, 2015).

The appearance of the UV absorption band centered at about 265 nm indicates the formation of tetragonal  $ZrO_2$  phase (Balaji *et al.*, 2017). By increasing the Y<sub>2</sub>O<sub>3</sub> content, the intensity of this band increases.

Villabona-Leal *et al.* (2014) consider that the UV-Vis absorption bands around 250 nm and 220 nm are due to different  $ZrO_2$  crystalline phases, namely the monoclinic (the 250 nm band) and tetragonal (the 220 nm) ones.

The UV-Vis feature at 350 nm is intense for the lowest content of  $Y_2O_3$ , flattened for the middle content and then disappears at the maximum yttrium content. This evolution can be assigned to the structural changes that occur in the samples. Thus, according to the XRD and FT-IR data, it is possible to associate the compositional evolution of this feature with the increase of the tetragonal ZrO<sub>2</sub> crystalline phase and the decrease of the monoclinic ZrO<sub>2</sub> crystalline phase, respectively.

Using the UV-Vis spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system, we determined the gap energy of the samples. The calculated gap energy values are presented in Fig. 4. The values for the direct gap energy are between 4.87 eV and 5.05 eV, whereas those for the indirect gap energy are between 5.12 and 5.14 eV. Both direct and indirect gap energy values decrease almost linearly with increasing the Y<sub>2</sub>O<sub>3</sub> content of the samples. Differences in the band gaps suggest different distributions and densities of the electronic states. The smaller values of the gap energy for the ceramic nanomixtures indicate smaller average crystallite sizes, in agreement with the SEM data.

#### 3.4. PL spectra

Pure zirconium and hafnium oxide samples containing only monoclinic phases show broadband PL emissions situated at about 490 nm. The microscopic origin of these PL bands is still disputable (Kiisk *et al.*, 2016).

The PL spectra of the  $3\text{HfO}_2 \cdot 15\text{SiO}_2 \cdot xY_2\text{O}_3 \cdot (82 - x)\text{ZrO}_2$ samples with x = 2, 7 and 12 mol%  $Y_2\text{O}_3$  are shown in Fig. 5. These spectra exhibit a broad band situated between 300 and 750 nm which is indicative of a white emission. This broad band consists of four PL bands of higher emission centered at about 335, 410, 450 and 470 nm and of two weak PL bands located at about 530 and 650 nm. The above-mentioned emission bands can be associated with structural defects such as oxygen vacancies in monoclinic  $\text{ZrO}_2$  and monoclinic HfO<sub>2</sub> structures.

The high-intensity PL band located at about 410 nm and the bands centered at about 530 and 650 nm can be caused by interstitial defects such as oxygen vacancies in  $HfO_2$  and  $ZrO_2$  crystalline phases with monoclinic structure when Hf and Zr atoms are seven-coordinated. The O atoms can occupy the threefold and fourfold positions and the average Hf—O bond length is 2.08 and 2.20 Å, respectively (Chuang *et al.*, 2012). Therefore, the fourfold-coordinated O atom around the four adjacent Hf atoms is more susceptible to forming a vacancy.



Figure 4

 $(\alpha h v)^{1/2}$  and  $(\alpha h v)^2$  as a function of photon energy for the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system; and the dependence of the optical band gap for the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system.

The stronger PL emission situated at  $\sim$  530 nm can be due to a fourfold-coordinated O atom in the HfO<sub>2</sub> structure.

As mentioned, the PL spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  samples are related to structural oxygen vacancies. The intensity of the PL emission bands increases with increasing  $Y_2O_3$  content of the samples. This behavior can be explained by considering that the small particle size was the main reason for the broad photoluminescence band.

#### 3.5. XANES data

XANES measurements in fluorescence mode were performed using the X-ray energy of the incident flux in the



Figure 5

Photoluminescence spectra of the samples in the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  composition where x = 2, 7 and  $12 \text{ mol}\% Y_2O_3$ .

range 9400–10400 eV. The  $L_3$ -edge of hafnium is located at about 9562 eV.

Zirconium and hafnium oxides are known to produce similar phase transitions under temperature and pressure variation. Under normal conditions of temperature and pressure, both oxides are in the monoclinic phase because the zirconium and hafnium cations have the same valence configuration with very similar ionic radii (0.78 Å for Zr<sup>+4</sup> and 0.76 Å for Hf<sup>+4</sup>). Since the crystallographic structures of the monoclinic phase of zirconium and hafnium are very close, many of the XRD peaks arise at nearly the same angles (Benyagoub, 2012). The transition temperature from the tetragonal to the monoclinic phase is higher in the case of  $HfO_2$  (~1700°C) than for  $ZrO_2$  (~1200°C). Note also that the  $Y^{+3}$  ions have radii (0.96 Å) larger than those of  $Zr^{+4}$  or  $Hf^{+4}$ ions and thus their insertion in the host matrix will produce lattice distortion and the elongation of the bonds with the oxygen ions.

Fig. 6 shows the XANES spectra for the Hf  $L_3$ -edge from 9525 to 9700 eV for the zirconate nanostructures in the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  composition with x = 2, 7 and 12 mol%  $Y_2O_3$ . The second derivative Hf  $L_3$ -edge XANES spectra are also presented in the lower graph of Fig. 6 and permit the variation in the energy to be distinguished.

The XANES features at the Hf  $L_3$ -edge of the studied composites show a strong main peak centered at about 9562 eV (feature A) corresponding to the Hf<sup>+4</sup> ions. In the post-edge region, a small peak located at about 9578 eV (feature B) and a broad peak situated around 9604 eV (feature C) were found.

The analysis of the XANES and second derivative spectra of the Hf  $L_3$ -edge indicates that the Y<sub>2</sub>O<sub>3</sub> content of samples



Hf  $L_3$ -edge XANES spectra of the studied samples in the 3HfO<sub>2</sub>·15SiO<sub>2</sub>·xY<sub>2</sub>O<sub>3</sub>·(82 - x)ZrO<sub>2</sub> system where x = 2, 7 and 12 mol% Y<sub>2</sub>O<sub>3</sub>. The lower graph shows the second derivative Hf  $L_3$ -edge XANES spectra.

do not affect the edge position of the peaks that remains the same for all the studied ceramics but influences the peaks' intensity. Thus, the intensity of the main peak (feature A) was the highest for sample x = 2, decreases for  $x = 7 \mod \% Y_2O_3$ and suffers a new increase for  $x = 12 \mod \% Y_2O_3$ . This variation of the peak intensity can be related to the different oxygen contents of the samples. The above-mentioned XANES feature (feature A) is located at about 9562 eV and corresponds to a dipole  $2p \rightarrow 5d$  transition due to the empty 5d states and the first EXAFS oscillation. The intensity of this feature depends on the filling of the *d* band of hafnium ions. In this view, note that the  $Hf^{+4}$  ions  $(5d^{0}6s^{0})$  have an empty 5d state while the  $Hf^{+3}$  ions  $(5d^{0}6s^{1})$  are partially filled. For the sample with  $x = 7 \mod \% Y_2O_3$ , the peak energy and intensity of the white lines are smaller than those in the case of the samples with x = 2 and 12 mol% Y<sub>2</sub>O<sub>3</sub>, suggesting a smaller number of 5d valences.

In conclusion, the Hf XANES data do not show significant differences with increasing  $Y_2O_3$  content of the samples suggesting an appreciable stability of the hafnium ions' +4 oxidation state and their microvicinity.

#### 3.6. EXAFS

To verify the origin of the XANES features at the Hf edge, the spectra corresponding to the monoclinic  $HfO_2$  structure were simulated.

The potential structural modifications and local structural changes caused to a host matrix by the addition of a dopant can be studied by EXAFS analysis and theoretical fitting procedures (Rada *et al.*, 2016). The EXAFS function was derived from the absorption data via the pre-edge and postedge subtraction. The magnitude of the Fourier transform of EXAFS oscillations at the Hf  $L_3$ -edge of the studied ceramics

Table 1

EXAFS structural parameters for hafnium in the studied samples.

*N*: model coordination number of the first coordinate sphere; *R*: the average Hf–O interatomic distance;  $\sigma^2$ : parameter of Debye–Waller thermal disorder (mean squared displacement);  $\Delta E$ : edge energy;  $\Delta R$ : shift of bonding distance between atoms.

Sample	Ν	<i>R</i> (Å)	$\sigma^2$	$\Delta E (eV)$	$\Delta R$ (Å)
$x = 2\% Y_2 O_3$	7	2.13041	0.00794	7.506	0.03041
$x = 7\% Y_2 O_3$	7	2.12626	0.00799	7.154	0.02626
$x = 12\% \mathbf{Y}_2 \mathbf{O}_3$	7	2.12059	0.00792	6.809	0.02059

is shown in Fig. 7. The EXAFS investigation detected three coordination shells around the hafnium ions for all the samples, namely Hf-O, Hf-Hf and a second Hf-Hf coordination shell at farther distances. The coordination numbers (N), average inter-atomic distances (R) and Debye-Waller disorder factors  $(\sigma^2)$  of the first coordination shell (Hf-O) around the hafnium ions are shown in Table 1.

The Fourier transform of the EXAFS signals at the Hf  $L_3$ edge of all samples shows two major oscillations. The first oscillation situated in the 0.8–2.2 Å range corresponds to the nearest neighbors around the hafnium and oxygen atoms. This main peak centered at about 1.7 Å indicates the distributions of oxygen atoms around hafnium atoms and is representative of the local structure order of Hf–O bonds in the first shell. The first coordination shell is almost uninfluenced by varying the Y<sub>2</sub>O<sub>3</sub> content from x = 2 to 7 mol% suggesting that the local order of Hf–O bonds is similar in both samples. After that, the increase of the Y<sub>2</sub>O<sub>3</sub> content up to x = 12 mol% leads to the increase of the peak related to the Hf–O bonds suggesting an increase of the average oxidation number of hafnium (Cho *et al.*, 2017).

The coordination number and the bond length of the first Hf-O shell are 8 and 2.22 Å, respectively, for the cubic  $HfO_2$  structure (Qi *et al.*, 2013). The EXAFS results presented in Table 1 show that in our samples the coordination numbers are around 7 and the bond lengths are ~2.12 Å that correspond to the monoclinic  $HfO_2$  structure. The lower coordination number and the bond length are due the charge-



Figure 7

Magnitude of the Hf  $L_3$ -edge Fourier transform of the EXAFS oscillation for the studied samples in the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$ composition where x = 2, 7 and  $12 \text{ mol}\% Y_2O_3$ .

compensating oxygen vacancies induced around the Hf atoms by addition of higher  $Y_2O_3$  contents.

The local structural parameters presented in Table 1 show a weak compositional variation. By increasing the  $Y_2O_3$  content of the samples the EXAFS first coordination shell distances (Hf-O distances) decrease from 2.13041 to 2.12059 Å. This evolution suggests that in sample  $x = 2 \mod^{9} Y_2O_3$  the local environment around the hafnium ions in the first coordination shell presents a lower order in comparison with the other samples.

The second major oscillation is situated in the 2.5-3.7 Å range and is related to the next nearest neighbors such as Hf-Hf or Hf-metals paths (Sharma et al., 2017). The other higher oscillations are due to the complex O-metal-O kinds of multiple scattering paths. Note that there is a significant difference in the second shell features. Thus, the fit of the second oscillation of the EXAFS signal shows that this oscillation consists of two path contributions with different lengths situated at  $\sim$ 2.9 and  $\sim$ 3.05 Å. The first contribution is related to the Hf-Zr (2.89 Å) and Hf-Hf (2.88 Å) bonding distances while the second contribution is associated with the Hf-Y bonding distance (3.06 Å). With increasing the yttrium content up to  $x = 12 \text{ mol}\% \text{ Y}_2\text{O}_3$ , the last peak gradually intensifies and shifts towards higher R values. This evolution depicts an increase in the Hf-metal bond length due to the enhancement of the crystallization process of the samples.

These compositional evolutions of the EXAFS data can be regarded as an indicator of the degree of crystallization at the hafnium local environment. For the sample with x = 1 mol% Y<sub>2</sub>O<sub>3</sub>, the contributions from external shells are more structured in comparison with those of the samples with x = 7 and 2 mol% Y<sub>2</sub>O<sub>3</sub>. The differences observed for the EXAFS

second coordination shell can be caused by the partial replacements of hafnium atoms by another metal atom as the second nearest neighbor of a hafnium atom, such as zirconium and yttrium atoms.

In conclusion, the EXAFS data suggest that the coordination number of the hafnium ions is around 7 and the Hf-Obond length is  $\sim 2.12$  Å, that are characteristic of the monoclinic HfO<sub>2</sub> structure. The contributions from the second coordination shells attributed to the Hf-metal paths become more structured by increasing the Y<sub>2</sub>O<sub>3</sub> content in the host ceramic matrix.

#### 3.7. SEM images

SEM micrographs of the samples the in  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  system with x = 2, 7 and  $12 \text{ mol}\% \text{ Y}_2\text{O}_3$  are shown in Fig. 8. Due to the increase of the  $Y_2O_3$  content, the samples  $x = 12 \mod \%$  will reach a more dense microstructure with smaller crystallite sizes. The highmagnification SEM image reveals that the structure of this sample consists of many domains with small and fine grains. The structural evolution of samples determined by increasing their Y2O3 content is in agreement with the XRD data. As was previously shown, the increase of the yttrium content determines the occurrence of yttrium-silicates precipitates (Y<sub>2</sub>SiO<sub>5</sub> and Y<sub>2</sub>Si<sub>2</sub>O<sub>7</sub> crystalline phases) that are distributed on the surface of zirconate samples.

#### 3.8. Raman spectroscopy

Characterization of zirconia systems by XRD may produce difficulties concerning the presence or absence of certain crystalline phases due to the highly broadened X-ray peaks.



Figure 8

SEM images of the studied samples in the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  composition with x = 2, 7 and  $12 \text{ mol}\% Y_2O_3$ .



Figure 9

Raman spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  samples with x = 2, 7 and  $12 \mod \% Y_2O_3$ .

The Raman spectroscopy technique can be useful in order to help the identification of zirconia polymorphs (Ghosh *et al.*, 2006; Hartmanova *et al.*, 1998).

The Raman spectra of the  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)$ ZrO<sub>2</sub> samples where x = 2, 7 and 12 mol% Y<sub>2</sub>O<sub>3</sub> are shown in Fig. 9. The Raman bands characteristic of the monoclinic ZrO<sub>2</sub> phase are those located at 180, 220, 225, 300, 335, 380, 470, 535, 555, 617 and 635 cm<sup>-1</sup> (Suya Prem Anand *et al.*, 2018). The tetragonal ZrO<sub>2</sub> phase is characterized by Raman bands centered at about 150, 255, 295, 306 and 430 cm<sup>-1</sup>. Thus, the Raman spectroscopic study of the studied zirconate samples offers a clear evidence of the presence of both monoclinic and tetragonal ZrO<sub>2</sub> structure in the prepared samples. The intensity of the Raman band centered at about 255 cm<sup>-1</sup> increases in the sample with x = 12 mol% Y<sub>2</sub>O<sub>3</sub> suggesting the monoclinic  $\rightarrow$  tetragonal phase transformation determined by the successive addition of Y<sub>2</sub>O<sub>3</sub> to the host matrix of samples.

The complex investigations of the ceramic nanomixtures (composites) were useful tools for distinguishing the zirconia and hafnia phases and the microstructural modifications at high dopant levels. The experiments reported in this paper demonstrate the structural modifications occurring in samples by doping.

The results show that yttrium as a dopant can be successfully used in order to modify structural and behavioral properties in a desired way. The studied materials can be used for optical applications after their gap band energies are lowered below 3 eV by appropriate procedures.

In order to show the applicability of these materials in biomedicine, investigations concerning their mechanical properties and chemical stability must complete the study.

#### 4. Conclusions

Zirconate nanostructured samples of  $3HfO_2 \cdot 15SiO_2 \cdot xY_2O_3 \cdot (82 - x)ZrO_2$  composition where x = 2, 7 and 12 mol% were prepared and studied by means of X-ray diffraction, X-ray

absorption spectroscopy, SEM analysis, photoluminescence, FT-IR, Raman and UV-Vis spectroscopy. XRD and XAS techniques were used to identify the crystalline phases from the studied samples. Thus, the presence of the following crystalline phases was observed:  $t-ZrO_2$ ,  $m-ZrO_2$ ,  $m-HfO_2$ ,  $Y_2SiO_5$  and  $Y_2SiO_5$ . SEM investigation highlighted the presence of  $Y_2SiO_5$  and  $Y_2Si_2O_7$  crystalline phases, in agreement with XRD data.

The photoluminescence emission bands are associated with structural defects such as the oxygen vacancies in monoclinic  $ZrO_2$  and monoclinic  $HfO_2$  structures. Based on the UV-Vis spectroscopy data, the gap energy of the samples was determined. The compositional evolution of gap energy values shows that the increase of the  $xY_2O_3$  content of the samples led to the enhancement of the crystallization process in agreement with the XRD data. Smaller values of the gap energy can also be related to smaller average crystallite sizes, in agreement with the SEM data.

FT-IR and Raman spectroscopies were useful tools for distinguishing the zirconia phases. Thus, it was observed that the increase of the  $Y_2O_3$  content of samples led to an increase of the t-ZrO<sub>2</sub> and silicate crystalline phases. The structural changes observed by the FT-IR and Raman investigations suggest that the yttrium ions play a network modifier role in the studied system.

The local structure around the Hf cations from the studied ceramic nanomixtures was studied by analysing the Hf  $L_3$ -edge EXAFS spectra obtained at room temperature. The EXAFS analysis suggests smaller deviations in bond length sets of Hf–O shells with increasing Y<sub>2</sub>O<sub>3</sub> content in the host ceramic matrix. The contributions from the second coordination shells attributed to the Hf–metal paths are more structured by increasing the Y<sub>2</sub>O<sub>3</sub> content in studied ceramics.

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#### References

- Balaji, S., Mandal, B. K., Ranjan, S., Dasgupta, N. & Chidambaram, R. (2017). J. Photochem. Photobiol. B, **170**, 125–133.
- Benyagoub, A. (2012). Acta Mater. 60, 5662-5669.
- Cho, D., Luebben, M., Wiefels, S., Lee, K. & Valov, I. (2017). Appl. Mater. Interfaces, 9, 19287–19295.
- Choi, D. & Choy, K. (2016). Electrochim. Acta, 218, 47-53.
- Chu, H., Hwang, W., Du, J., Chen, K. & Wang, M. (2016). J. Alloys Compd. 678, 518–526.
- Chuang, S., Lin, H. & Chen, C. (2012). J. Alloys Compd. 534, 42-46.
- Dehelean, A., Rada, S., Popa, A., Suciu, R. C. & Culea, E. (2016). J. Lumin. 177, 65–70.
- Fan, Y., Cheng, S., Wang, H., Tian, J., Xie, S., Pei, Y., Qiao, M. & Zong, B. (2017). Appl. Catal. Environ. 217, 331–341.
- Ferdov, S., Ferreira, R. A. S., Lin, Z. & Wu, Z. (2012). J. Solid State Chem. 190, 18–23.
- Ghosh, A., Suri, A. K., Pandey, M., Thomas, S., Rama Mohan, T. R. & Rao, B. T. (2006). *Mater. Lett.* **60**, 1170–1173.
- Glazneva, T. S., Kaichev, V. V., Paukshtis, E. A., Khabibulin, D. F., Lapina, O. B., Bal'zhinimaev, B. S., Zhurba, E. N., Lavrinovich, I. A.,

## research papers

Gavrikova, I. N., Shumskii, V. I. & Trofimov, A. N. (2012). J. Non-Cryst. Solids, 358, 1053–1058.

- Hartmanova, I., Thurzo, M., Jergel, I., Bartos, F., Kadlec, V., Zelezny, D., Tunega, F., Kundracik, F., Chromik, S. & Brunel, M. (1998). *J. Mater. Sci.* 33, 969–975.
- Karunarathne, M. C., Baumann, J. W., Heeg, M. J., Martin, P. D. & Winter, C. H. (2017). J. Organomet. Chem. 847, 204–212.
- Khomenkova, L., An, Y., Khomenkov, D., Portier, X., Labbé, C. & Gourbilleau, F. (2014). *Physica B*, 453, 100–106.
- Kiisk, V., Puust, L., Utt, K., Maaroos, A., Mändar, H., Viviani, E., Piccinelli, F., Saar, R., Joost, U. & Sildos, I. (2016). *J. Lumin.* **174**, 49–55.
- Kwon, O. H., Jang, C., Lee, J., Jeong, H. Y., Kwon, Y., Joo, J. H. & Kim, H. (2017). *Ceram. Int.* 43, 8236–8245.
- Lee, S. W. & Condrate, R. A. (1988). J. Mater. Sci. 23, 2951-2959.
- Moll, S., Thomé, L., Sattonnay, G., Garrido, F., Vincent, L., Jagielski, J. & Thomé, T. (2009). Vacuum, 83, S61–S64.
- Najafinezhad, A., Abdellahi, M., Ghayour, H., Soheily, A., Chami, A. & Khandan, A. (2017). *Mater. Sci. Eng. C*, **72**, 259–267.
- Ortiz, A., Alonso, J. C. & Haro-Poniatowski, E. (2005). *J. Elec Mater.* 34, 150–155.
- Ponnilavan, V., Vasanthavel, S., Singh, R. K. & Kannan, S. (2015). *Ceram. Int.* 41, 7632–7639.
- Qi, J. & Zhou, X. (2015). Colloids Surf. A Physicochem. Eng. Asp. 487, 26–34.
- Qi, Z., Cheng, X., Zhang, G., Li, T., Wang, Y., Shao, T., Li, C. & He, B. (2013). Mater. Res. Bull. 48, 2720–2723.

- Qin, H., Liu, H., Sang, Y., Lv, Y., Zhang, X., Zhang, Y., Ohachi, T. & Wang, J. (2012). CrystEngComm, 14, 1783–1789.
- Rada, M., Aldea, N., Wu, Z. H., Jing, Z., Rada, S., Culea, E., Macavei, S., Balan, R., Suciu, R. C., Erhan, R. V. & Bodnarchuk, V. (2016). *Journal of Non-Crystalline Solids*, 437, 10–16.
- Rada, M., Zagrai, M., Rada, S., Bot, A. & Culea, E. (2017). J. Alloys Compd. 705, 327–332.
- Rada, M., Zhang, J., Zhongua, W., Zagrai, M., Magerusan, L. & Rada, S. (2018). *Anal. Lett.* **51**, 2544–2550.
- Rada, S., Culea, E. & Rada, M. (2018). Chem. Phys. Lett. 696, 92-99.
- Rada, S., Rada, M. & Culea, E. (2013). J. Alloys Compd. 552, 10-13.
- Rada, S., Rada, M., Erhan, R. V., Bodnarchuk, V., Barbu-Tudoran, L. & Culea, E. (2019). J. Alloys Compd. 770, 395–404.
- Sharma, A., Varshney, M., Shin, H., Chae, K. & Won, S. O. (2017). Spectrochim. Acta A Mol. Biomol. Spectrosc. 173, 549–555.
- Simhan, R. G. (1983). J. Non-Cryst. Solids, 54, 335-343.
- Singh, S., Kalia, G. & Singh, K. (2015). J. Mol. Struct. 1086, 239-245.
- Suya Prem Anand, P., Arunachalam, N. & Vijayaraghavan, L. (2018). J. Mech. Behav. Biomed. Mater. 86, 122–130.
- Torres, D. I. & Llopis, J. (2009). Superlattice Microstruct. 45, 482-488.
- Villabona-Leal, E. G., Diaz-Torres, L. A., Desirena, H., Rodríguez-López, J. L., Pérez, E. & Meza, O. (2014). J. Lumin. 146, 398–403.
- Vitanov, P., Harizanova, A. & Ivanova, T. (2014). J. Phys. Conf. Ser. 514, 012011.
- Wu, L., Wu, C. G. & Chang, Y. (1992). J. Electron. Mater. 21, 217-222.
- Zhen, Q., Chen, R., Van, K. & Li, R. (2007). J. Rare Earths, 25, 199–203.